

Time reversal (TR)

$$f(t) \rightarrow f(-t) \rightarrow \text{time reversed counterpart.} \quad t' = -t$$

$$\text{Classically: } v = \frac{dx(t)}{dt} \rightarrow \frac{dx(-t)}{dt} = \frac{dx(t')}{dt'} \frac{dt'}{dt} = v(-1) = -v$$

Given $\{q_i, v_i\}$, its TR counterpart is $\{q_i, -v_i\}$

If energy is same then H must have TR symmetry.
(For example if H does not depend on odd powers of \dot{q}_i)

Note that you can not go from one description to the other description unless you flip the force but

$$\text{TR does not flip force: } \frac{d^2x}{dt^2} \xrightarrow{\text{TR}} \frac{d^2x}{dt^2}$$

So there are discrete snap shots.

TR symmetry \rightarrow discrete symmetry.

$$\text{In QM: motion} \rightarrow \frac{1}{m} \langle \psi | \hat{p} | \psi \rangle$$

$i \dot{q}$ $i \dot{q}$

$$\text{let } \psi = c_2 e^{i q x} + c_2' e^{i q' x}$$

$$\langle \psi | \hat{p} | \psi \rangle = |c_2|^2 q + |c_2'|^2 q'$$

$$\psi^* = c_2^* e^{-i q x} + c_2'^* e^{-i q' x}$$

$$\langle \psi^* | \hat{p} | \psi^* \rangle = |c_2|^2 (-q) + |c_2'|^2 (-q')$$

$\therefore \langle \psi | \hat{p} | \psi \rangle$ and $\langle \psi^* | \hat{p} | \psi^* \rangle$
are motion reversed!

\therefore In QM: motion \rightarrow Complex conjugation of space part.
reversed
We can see it more formally by coming through TDSE.

$$\text{TDSE: } i \hbar \frac{\partial}{\partial t} \psi(x,t) = H \psi(x,t)$$

$$\Rightarrow \langle \psi | (i \hbar \frac{\partial}{\partial t}) | \psi \rangle = E \text{ Real } \text{--- (A)}$$

$$\dots \dots \dots \langle \psi(-t) \rangle = E$$

$$TR: t \rightarrow -t \quad \langle \psi(-t) | i\hbar \frac{\partial}{\partial t} | \psi(t) \rangle = E \quad \text{--- (B)}$$

$$[\text{or } -t = t']$$

$$\frac{\partial t'}{\partial t} \rightarrow \bar{} \quad \langle \psi(t') | i\hbar \frac{\partial}{\partial t'} | \psi(t') \rangle = E \quad \text{--- (B)}$$

t' can as well be set to t :

$$\therefore - \langle \psi(t) | i\hbar \frac{\partial}{\partial t} | \psi(t) \rangle = E \quad \text{contradicts (A)!}$$

However if we do complex conjugation of (B)

$$[t' = -t] \quad \langle \psi^*(-t) | -i\hbar \frac{\partial}{\partial t} | \psi^*(-t) \rangle = E$$

$$\Rightarrow \quad \langle \psi^*(t') | i\hbar \frac{\partial}{\partial t'} | \psi^*(t') \rangle = E$$

No contradiction with (A) 😊

" In QM \xrightarrow{TR} (set t to $-t$) + Complex conjugation.

$$\Rightarrow \psi(x, t) \xrightarrow{TR} \psi^*(x, -t)$$

Note that we do this only in wave function.

NOT in operators. Because they are our measuring tools, and $\Psi(x,t)$ and $\Psi^*(x,-t)$ are two states of the same system. They could as well be the same state. Think of $e^{i\epsilon x}$ with $\epsilon=0$.

Define TR operator as $\hat{\theta}$ which does complex conjugation and converts t to $-t$ at all states.

Let us see how $\hat{\theta}$ operates on Bloch states.

$$\hat{\theta} \hat{H} \Psi_{\epsilon n} = E_{\epsilon n} \hat{\theta} \Psi_{\epsilon n} \quad \text{since } E_{\epsilon n} \text{ is real.}$$

$$\Rightarrow \hat{H} \hat{\theta} \Psi_{\epsilon n} = E_{\epsilon n} \hat{\theta} \Psi_{\epsilon n} \quad \text{since } \hat{H} \text{ is real.}$$

$$\Rightarrow [\hat{H} \hat{\theta}] = 0 \Rightarrow \Psi_{\epsilon n} \text{ and } \hat{\theta} \Psi_{\epsilon n} \text{ are degenerate!}$$

What is $\hat{\theta} \Psi_{\epsilon n}$?

$$\hat{\theta} \Psi_{\epsilon n} = \Psi_{-\epsilon n}$$

$$\theta \psi_{\mathbf{q}_n} = \theta e^{-i\mathbf{q}_n \cdot \mathbf{r}} u_{\mathbf{q}_n}(\mathbf{r})$$

$$= e^{-i\mathbf{q}_n \cdot \mathbf{r}} \hat{\theta} u_{\mathbf{q}_n}(\mathbf{r})$$

Remember,

$$\hat{H}_{\mathbf{q}} u_{\mathbf{q}} = E_{\mathbf{q}} u_{\mathbf{q}}$$

$$\hat{\theta} \hat{H}_{\mathbf{q}} u_{\mathbf{q}} = E_{\mathbf{q}} \hat{\theta} u_{\mathbf{q}}$$

$$\hat{\theta} \hat{H}_{\mathbf{q}} u_{\mathbf{q}_n} = \hat{\theta} \left[\frac{(\mathbf{p} + \hbar \mathbf{q})^2}{2m} + v \right] u_{\mathbf{q}_n}$$

$$= \hat{\theta} \frac{1}{2m} \left[\hbar^2 \mathbf{q}^2 - \hbar^2 \frac{\partial^2}{\partial \mathbf{r}^2} + 2i\hbar^2 \mathbf{q} \cdot \frac{\partial}{\partial \mathbf{r}} + v \right] u_{\mathbf{q}_n}^*$$

$$= \frac{1}{2m} \left[\hbar^2 \mathbf{q}^2 - \hbar^2 \frac{\partial^2}{\partial \mathbf{r}^2} - 2i\hbar^2 \mathbf{q} \cdot \frac{\partial}{\partial \mathbf{r}} + v \right] u_{\mathbf{q}_n}^*$$

$$= \left[\frac{(\mathbf{p} - \hbar \mathbf{q})^2}{2m} + v \right] \hat{\theta} u_{\mathbf{q}_n} \quad \hookrightarrow \hat{\theta} u_{\mathbf{q}_n}$$

$$\therefore \hat{\theta}(\hat{H}_q \psi_q) = E_{qn} \hat{\theta} \psi_q$$

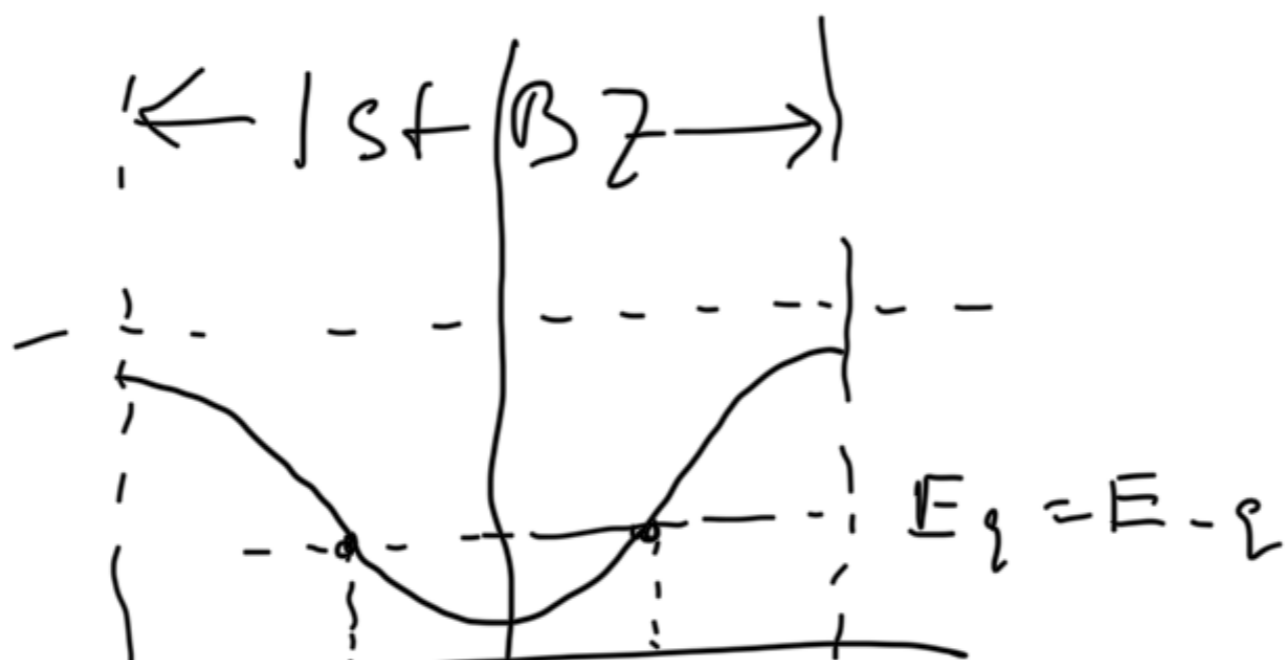
$$\Rightarrow \hat{H}_{-q} \hat{\theta} \psi_{qn} = E_{qn} \hat{\theta} \psi_{qn}$$

Compare with $\hat{H}_{-q} \psi_{-qn} = E_{-qn} \psi_{-qn}$

$$\Rightarrow \hat{\theta} \psi_{qn} = \psi_{-qn} \text{ and } E_{qn} = E_{-qn}$$

$$\therefore \hat{\theta} \psi_{qn} = e^{-i\varphi_n} \hat{\theta} \psi_{qn} = e^{-i\varphi_n} \psi_{-qn} = \psi_{-qn}$$

$$\therefore [\hat{H}, \hat{\theta}] = 0 \Rightarrow \hat{\theta} \psi_{qn} = \psi_{-qn} \Rightarrow \psi_{qn} \text{ and } \psi_{-qn} \text{ degenerate!}$$



$$\begin{array}{c|c} \hline & -q \\ \hline -\frac{\Delta}{a} & \frac{\Delta}{a} \\ \hline \end{array}$$

←→

Irreducible BZ

If $[H\hat{Q}] = 0$, it is sufficient to know $\{E_{2n}, \psi_{2n}\}$ of both
of the 1st BZ.